


$$\hat{H} \Psi_n(r) = E_n \Psi_n(r)$$

Remember: This is the space part. The time dependent part is $e^{-i\omega_n t}$, where $E_n = \hbar \omega_n$:

$$\Psi_n(r, t) = \Psi_n(r) e^{-i\omega_n t}$$

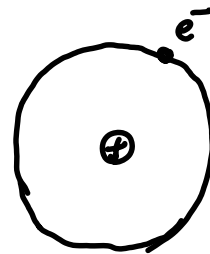
Example: what is the \hat{H} for these potentials:

1D Simple Harmonic Oscillator (SHO),

$$\hat{H} = \underbrace{\frac{-\hbar^2}{2m} \frac{d^2}{dx^2}}_T + \underbrace{\frac{1}{2} m \omega^2 x^2}_V$$


Hydrogen atom,

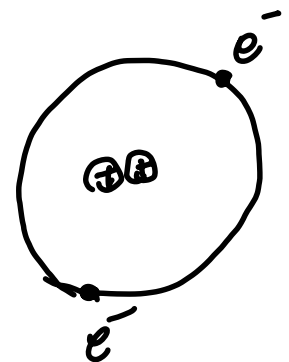
$$\hat{H} = -\underbrace{\frac{\hbar^2}{2m} \nabla^2}_{T_e} - \underbrace{\frac{\hbar^2}{2M} \nabla^2}_{T_p} - \underbrace{\frac{e^2}{4\pi\epsilon_0 r}}_{V_{ep}}$$



Helium atom,

$$\hat{H} = T_{e1} + T_{e2} + T_{p1} + T_{p2} +$$

$$V_{e1e2} + V_{e1p1} + V_{e1p2} + V_{e2p1} + V_{e2p2} + V_{p1p2}$$



Using the Schrödinger wave equation

$$\hat{H} \Psi_n(r) = \left(-\frac{\hbar^2}{2m} \nabla^2 + V(r) \right) \Psi_n(r) = E_n \Psi_n(r)$$

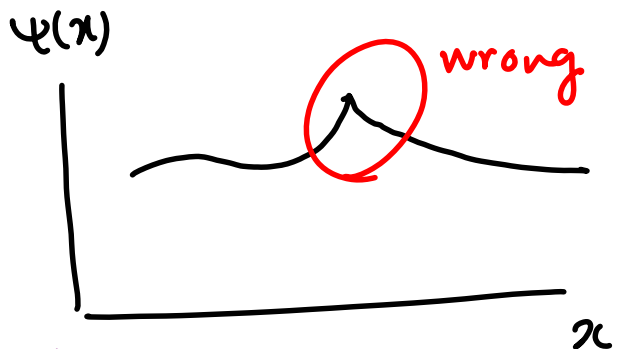
$$\Psi_n(r, t) = \underbrace{\Psi_n(r)}_{\text{Stationary state (time independent)}} e^{-i\omega_n t} \quad E_n = \hbar\omega_n$$

Boundary conditions:

① Ψ is continuous



② $\frac{d\Psi}{dx}$ is continuous if V is finite.



③ $\Psi(x \rightarrow \pm\infty) = 0$

wave function normalization

$$\int_{-\infty}^{\infty} |\Psi(r)|^2 d^3r = 1$$

Wavefunctions are orthonormal

$$\hat{H} \Psi_n = E_n \Psi_n \quad \Rightarrow \quad \Psi_1, \Psi_2, \dots$$

$$\int_{-\infty}^{\infty} \Psi_n^*(r) \Psi_m(r) d^3r = \delta_{nm} \begin{cases} 1 & n=m \\ 0 & n \neq m \end{cases}$$

↓
Kronecker delta

Completeness

Ψ_n 's make a complete space.

So any $\Psi(r)$ can be expanded

by Ψ_n 's:

$$\Psi(r) = \sum_n a_n \Psi_n(r)$$

Inversion Symmetry in $V(r)$

$$\underbrace{\left(-\frac{\hbar^2}{2m} \nabla^2 + V(r) \right)}_{\hat{H}_n} \Psi_n = E_n \Psi_n$$

If $V(r)$ is symmetric: $V(r) = V(-r)$,
eigenfunctions have either odd or even
parity:

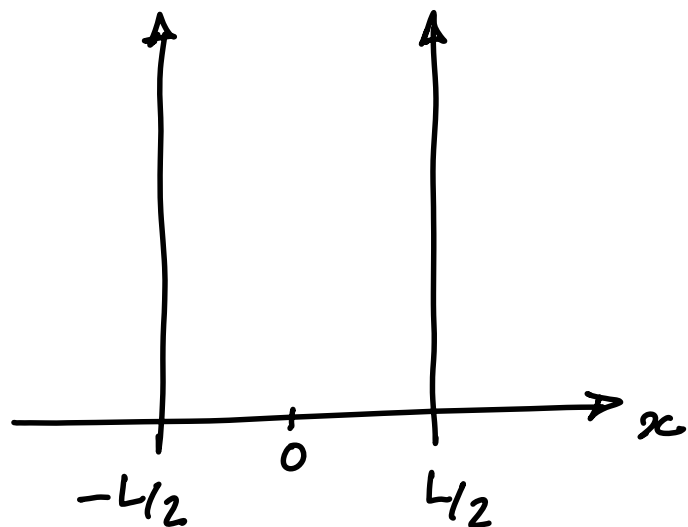
Even parity: $\Psi_n(r) = \Psi_n(-r)$

Odd parity: $\Psi_n(r) = -\Psi_n(-r)$

Example 1D rectangular potential well
with infinite barrier energy

$$V=0 \quad -\frac{L}{2} < x < \frac{L}{2}$$

$$V=\infty \quad \text{elsewhere}$$



$$-\frac{\hbar^2}{2m} \frac{d^2 \Psi}{dx^2} = E \Psi$$

$$\Psi'' = - \underbrace{\frac{2mE}{\hbar^2}}_{k^2} \Psi$$

$$\Rightarrow \psi = A e^{ikx} + B e^{-ikx}$$

B.C.

$$\psi(-\frac{L}{2}) = \psi(\frac{L}{2}) = 0 \Rightarrow$$

$$\left\{ \begin{array}{l} \psi(-\frac{L}{2}) = A e^{-ikL/2} + B e^{ikL/2} = 0 \\ \psi(\frac{L}{2}) = A e^{ikL/2} + B e^{-ikL/2} = 0 \end{array} \right.$$

$$\left| \begin{array}{cc} e^{-ikL/2} & e^{ikL/2} \\ e^{ikL/2} & e^{-ikL/2} \end{array} \right| = \underbrace{e^{-ikL} - e^{ikL}} = 0$$

↓

$$\cancel{\cos kL} - i \sin kL - (\cancel{\cos kL} + i \sin kL)$$

$$\Rightarrow \sin kL = 0 \Rightarrow k_n L = n\pi \quad n=1, 2, \dots$$

$$\Rightarrow \boxed{\begin{array}{l} k_n = \frac{n\pi}{L} \\ E = \frac{\hbar^2 k^2}{2m} \Rightarrow E_n = \frac{\hbar^2}{2m} \frac{n^2 \pi^2}{L^2} \end{array}}$$

$$\left. \begin{array}{l} k_n = \frac{n\pi}{L} \\ k = \frac{2\pi}{\lambda} \end{array} \right\} \frac{2\pi}{\lambda_n} = \frac{n\pi}{L} \Rightarrow \boxed{L = n \frac{\lambda_n}{2}} \quad n=1, 2, \dots$$

